

***Ab initio* molecular dynamics with simultaneous electron and phonon excitations: application to the relaxation of hot atoms and molecules on metal surfaces**

D. Novko,¹ M. Blanco-Rey,^{2,1} J.I. Juaristi,^{2,3,1} and M. Alducin^{3,1}

¹ *Donostia International Physics Center (DIPC), Paseo Manuel de Lardizabal 4, 20018 Donostia-San Sebastián, Spain*

² *Departamento de Física de Materiales, Facultad de Químicas UPV/EHU, Apartado 1072, 20080 Donostia-San Sebastián, Spain*

³ *Centro de Física de Materiales CFM/MPC (CSIC-UPV/EHU), Paseo Manuel de Lardizabal 5, 20018 Donostia-San Sebastián, Spain*

(Dated: July 28, 2016)

The relaxation dynamics of hot H, N and N₂ on Pd(100), Ag(111) and Fe(110), respectively, is studied by means of *ab initio* molecular dynamics with electronic friction (AIMDEF). This method is adapted here to account for the electron density changes caused by lattice vibrations, thus treating on an equal footing electron-hole (e-h) pair and phonon excitations. We find that even if the latter increasingly dominate the heavier the hot species, the contribution of e-h pairs is by no means negligible in these cases because it gains relevance at the last stage of the relaxation process. The quantitative details of energy dissipation depend on the interplay of the potential energy surface, electronic structure and kinetic factors.

PACS numbers: 34.35.+a, 68.43.-h, 34.50.Bw

In dynamic gas-surface environments, where gas-phase atomic and molecular species impinge on the surface at energies of the order of up to a few eV, energy dissipation occurs by excitation of electron-hole (e-h) pairs and excitation of lattice vibrations, i.e. phonons [1–17]. In the adsorption processes of atomic and molecular species, dissociative as well as non-dissociative, the species trapped by the surface gradually lose their energy until they become thermalized on the surface. The competition between the e-h pairs and phonons channels governs the relaxation dynamics of the transient *hot* species, and thus it plays a decisive role in the system reactivity properties. The reason is that it rules the traveled length and relaxation time of a hot atom or molecule on the surface and, consequently, the probability to undergo a recombination reaction with another adsorbate [18–23].

Recent *ab initio* molecular dynamics (AIMD) simulations with electronic friction (AIMDEF) have shown that e-h pair excitations are the dominant relaxation mechanism for hot H atoms on Pd(100) that originate from the dissociative adsorption of H₂ [16]. More particularly, this channel dissipates energy at a five times faster rate than the phonons channel [10]. The two main reasons behind this behavior are the long H-Pd interaction time, of hundreds of fs, and the low adsorbate-to-surface atom mass ratio, $\gamma = m_{\text{H}}/m_{\text{Pd}} = 0.0094$. The case of H on Pd(100) represents a limiting case. For heavier adsorbates, the relative weight of e-h pairs and phonons in the energy loss is expected to vary. The energy transfer to the substrate will be determined not only by kinetic factors, such as the value of γ and the incidence conditions, but also by the topography of the multidimensional potential energy surface (PES) and the electronic structure details of the configurations probed along the relaxation trajectory.

In this paper, we investigate the relaxation dynamics of

hot species in three adsorption scenarios representative of different energy loss regimes. We have chosen atomic N on Ag(111) ($\gamma = 0.13$), N₂ on Fe(110) ($\gamma = 0.5$) and the aforementioned H on Pd(100) as case studies. Our choice is also motivated by the results reported from chemi-current experiments showing that the number of low-energy metal electrons excited during the adsorption of different gas species, i.e., the chemi-current intensity, scales with the adsorption energy [4, 5]. The wide range of E_{ads} values covered by the present case studies, between 0.2 and 2.3 eV, allows to elucidate the observed scaling law. Modeling on an equal footing the e-h pair and phonon dissipation mechanisms is crucial to properly account for the hot species relaxation dynamics, irrespective of the considered kinetic factors. Thus, our challenge here has been to adapt the efficient AIMDEF methodology in order to accurately treat e-h pair excitations in those situations of non-negligible electron density changes produced by the surface atoms movement.

In AIMDEF [16, 24], the effect of low energy e-h pair excitations on the gas species dynamics is included by a friction force that is calculated *on-the-fly* within the local density friction approximation (LDFA) [25]. This approximation assumes that the friction coefficient $\eta(\mathbf{r}_A)$ acting on each atom of the hot species at its position \mathbf{r}_A is the same friction coefficient the atom would experience when embedded in a homogeneous free electron gas with electron density $n(\mathbf{r}_A)$. The latter is the density of the *bare* surface at \mathbf{r}_A , which is calculated with density functional theory (DFT). Obviously, obtaining $n(\mathbf{r}_A)$ *on-the-fly* becomes a complicated task in simulations where surface atoms are allowed to move. The problem is that at each time step of the trajectory $n(\mathbf{r}_A)$ should be that of the actual configuration of the (moving) surface atoms at this precise instant. Here, we propose

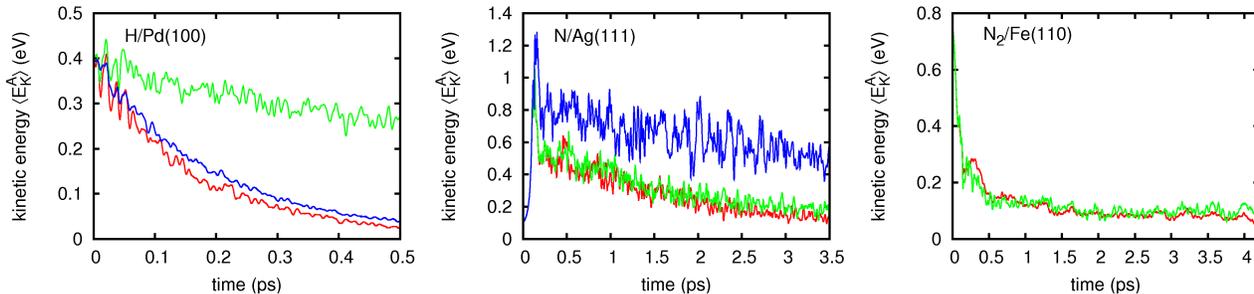


FIG. 1. Adsorbate kinetic energies averaged over AIMDEF-simulated trajectories of H on Pd(100), N on Ag(111) and N_2 on Fe(110), respectively. The blue (red) lines correspond to FS+EF (NFS+EF) calculations, and the green lines correspond to NFS calculations.

a new scheme to evaluate the embedding electronic density that allows to incorporate density variations due to surface atoms movement. The idea consists in evaluating the gas-species electron density within a Hirshfeld partitioning analysis [26] and use it to remove its contribution from the self-consistent gas-surface electron density $n^{\text{SCF}}(\mathbf{r}_A)$. Therefore, the embedding bare surface density for each atom can be calculated *on-the-fly* at each integration step as,

$$n_{\text{sur}}(\mathbf{r}_A) = n^{\text{SCF}}(\mathbf{r}_A) \left[1 - \frac{\sum_j^{N'} n_j^{\text{atom}}(\mathbf{r}_A)}{\sum_i^N n_i^{\text{atom}}(\mathbf{r}_A)} \right], \quad (1)$$

where the summations on i and j run over all the N atoms of the system and over the N' atoms conforming the gas-species, respectively, and $n_{i,j}^{\text{atom}}$ is the electron density associated to the isolated atoms. The Hirshfeld partitioning scheme has successfully been applied to study vibrational lifetimes of molecular adsorbates within the LDFA framework [27]. Here we have checked that the energy losses obtained in frozen-surface simulations by using either the self-consistent bare surface electron density or the present n_{sur} are basically undistinguishable.

We start revisiting the relaxation of hot H atoms formed from H_2 dissociation on Pd(100) at normal incidence and initial kinetic energy $E_i(H_2) = 0.5$ eV. Our simulations (178 individual H trajectories) start at the instant ($t = 0$) in which H_2 is already dissociated (see Ref. [16] for further technical details). The adsorption cases of N atoms on Ag(111) and N_2 on Fe(110) are studied for normal incidence and initial energies $E_i = 0.1$ and 0.75 eV, respectively. These values are chosen to ensure adsorption probabilities $S_0 > 0.85$ [28] and > 0.7 [29], respectively, that will yield reliable statistical averages. A total of 20(80) trajectories departing from height $z_i = 4$ Å are simulated using as time step 0.5(0.7) fs for N atoms (N_2 molecules). Special care is taken in describing properly the N spin-state at different heights. The supercell size used in each system ensures negligible lateral interactions [30]. In order to disentangle the contributions of e-h pair and phonon excitations we perform three types

of AIMD simulations on each system: (i) frozen-surface with electronic friction (FS+EF), (ii) non frozen-surface without electronic friction (NFS), where the two outer metal layers are allowed to move, and (iii) non frozen-surface with electronic friction (NFS+EF).

Figure 1 shows, for each simulation type and system, the adsorbate kinetic energy as a function of time averaged over the available number of adsorption trajectories, $\langle E_K^A \rangle$. Starting with the H/Pd(100) case, we first note that the NFS+EF result is not significantly affected by the use of n_{sur} instead of the bare frozen surface density approach taken in Ref. [16]. The reason is that Pd atoms are only slightly displaced by collisions with H. The figure shows the two typical features of a system in which relaxation is dominated by e-h pair excitations: (i) $\langle E_K^A \rangle(t)$ decays more rapidly in the FS+EF than in the NFS simulations and, as a consequence, (ii) the NFS+EF results, which include both dissipation channels, lie close to the FS+EF curve, which only includes the dominant electronic channel.

In contrast, the N/Ag(111) results are those representative of a system in which relaxation is dominated by phonon excitations, i.e., a much faster decay of $\langle E_K^A \rangle(t)$ when surface atom movement is allowed, very similar irrespectively of whether electronic friction is included (NFS+EF) or not (NFS). Regarding the details of the dynamics, at the beginning of the N/Ag(111) simulation ($t < 0.2$ ps), we observe a rapid kinetic energy gain, which is the result of the barrierless strong attraction felt by the N atom towards the Ag(111) surface. This increase is followed by a rapid decrease that is a consequence of N probing the repulsive part of the PES. In the NFS and NFS+EF cases there is an additional energy transfer to the surface atoms acting as soon as N approaches the surface, as we will discuss below. Therefore, after the first collision the N atom is left with about 0.5 eV of kinetic energy (0.8 eV in the FS+EF case) to be lost as a regular hot atom.

The case of nondissociative adsorption of N_2 on Fe(110) is even more extreme regarding the dominant role played by phonon over e-h pair excitations. In fact,

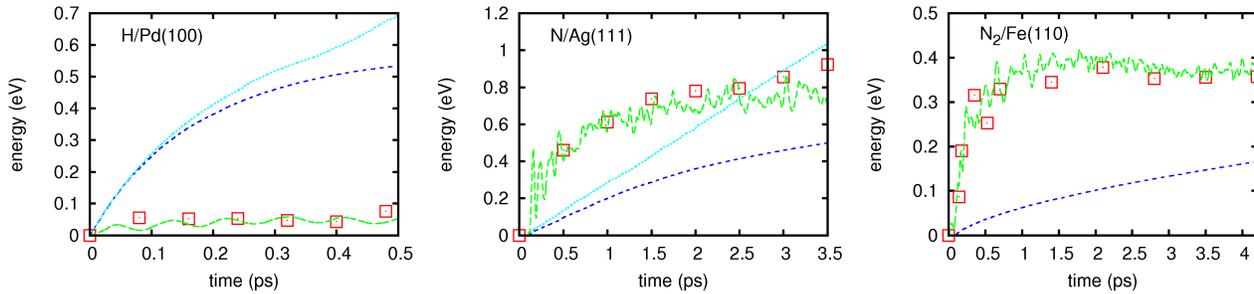


FIG. 2. Contributions to the energy transferred to the surface upon adsorption of H on Pd(100), N on Ag(111) and N_2 on Fe(110) in NFS+EF simulations: energy lost into e-h pair excitations, $\langle E_{\text{ehp}} \rangle$ (dark blue), kinetic energy of surface atoms, $\langle E_K^S \rangle$ (green) and potential energy of the surface, $\langle E_P^S \rangle$ (red squares). $\langle E_{\text{ehp}} \rangle$ for FS+EF simulations is represented in light blue.

there are no molecular adsorption events in the absence of surface atom movement. Furthermore, the NFS and NFS+EF results are almost undistinguishable. In contrast to N/Ag(111), we do not observe the fast kinetic energy gain upon approach of the projectile to the surface. The reason is the presence of energy barriers at $z \sim 2.5 \text{ \AA}$ [29, 31] that the molecules have to overcome before accessing the adsorption wells. As a result, we observe a drop of $\langle E_K^A \rangle$ starting at $t = 0$ that is the combined effect of an increase of the potential energy and of energy transfer to the surface lattice. Moreover, unlike in the two previously studied atomic adsorption cases, N_2 molecules undergo little lateral displacement as hot species on the surface.

All in all, Fig. 1 highlights that, under the usual conditions relevant for gas-surface reactions, e-h pair excitations will dominate energy dissipation only for light atoms ($\gamma \ll 1$). In principle, the kinetic energy of a free atom in a free electron gas (FEG) decays at a rate $2\eta/m_A$, where m_A is the atom mass and its friction coefficient η depends non-trivially on the atomic number Z and the FEG density (Z -oscillations) [32]. Nevertheless, for the typical electron densities probed by hot atoms and molecules on metal surfaces, η varies slowly with Z [32–34] and thus the electronic decay rate is dominated by m_A . As a general trend, this causes the behavior observed in Fig. 1, although the actual decay rates deviate from the free-atom values due to the PES topography of each system [16].

Next, we analyze the amount of energy transferred from the adsorbates to e-h pairs and to the lattice. For a single atomic trajectory, the e-h pair energy contribution as a function of time can be evaluated by integration of the friction force on each gas atom,

$$E_{\text{ehp}}(t) = \int_0^t \eta(\mathbf{r}_A(t')) |\mathbf{v}_A(t')|^2 dt', \quad (2)$$

where $\mathbf{v}_A(t')$ is the instantaneous atom velocity. The average value of this quantity over the available number of trajectories $\langle E_{\text{ehp}}(t) \rangle$ is shown in Fig. 2 for the FS+EF and NFS+EF simulations performed on the three sys-

tems (for N_2 the contributions of both atoms are added). In order to extract the energy transferred to the lattice atoms, we need both the instantaneous kinetic energy of the surface atoms and the instantaneous variation of the potential energy due to lattice distortions. The former is provided *on-the-fly* in the NFS+EF simulations and its statistical average over all the trajectories $\langle E_K^S(t) \rangle$ is also shown in Fig. 2. The energy stored as potential energy of the surface atoms $\langle E_P^S \rangle$ can be obtained from an analysis *a posteriori* of the simulated trajectories. For a few individual time steps $\{t_n\}$, we remove the adsorbate and consider only the surface atomic coordinates. In these distorted surface configurations, we evaluate the increase in the potential energy with respect to the equilibrium bare-surface configuration. This calculation is carried out for the available trajectories and then averaged to obtain estimates of $\langle E_P^S(t_n) \rangle$. Figure 2 shows for the three systems that $\langle E_P^S(t_n) \rangle$ and $\langle E_K^S(t) \rangle$ lie close to each other. This can be interpreted as a fingerprint of a harmonic oscillatory regime for the surface atoms. The average total energy transferred to the lattice, i.e., the phononic part of the energy loss, is $\langle E_{\text{ph}} \rangle = \langle E_K^S \rangle + \langle E_P^S \rangle$.

In the e-h pair dominated H/Pd(110) system, Fig. 2 shows that after just 0.5 ps the energy transferred to e-h pairs is around five times larger than the total energy transferred to the lattice ($\langle E_{\text{ph}} \rangle \simeq 0.1 \text{ eV}$). We also observe that energy dissipation to each of the channels is not additive, since inclusion of phonon excitations reduces the amount of energy going to e-h pairs. This effect is more pronounced in the N/Ag(111) system, in which the energy transfer to e-h pair excitations is around a factor two larger in the FS+EF than in the NFS+EF simulations after 3.5 ps. The main reason is that the average \mathbf{v}_A values entering Eq. (2) are lower in the NFS+EF than in the FS+EF simulations due to the additional energy release into the competing phononic channel.

There are some common features in N/Ag(111) and N_2 /Fe(110), i.e., the systems in which phonons dominate the adsorption relaxation. In the first stages of the dynamics, at $t < 0.5 \text{ ps}$ in N_2 /Fe(110) and $t < 1.5 \text{ ps}$ in N/Ag(111), a rapid energy transfer takes place from

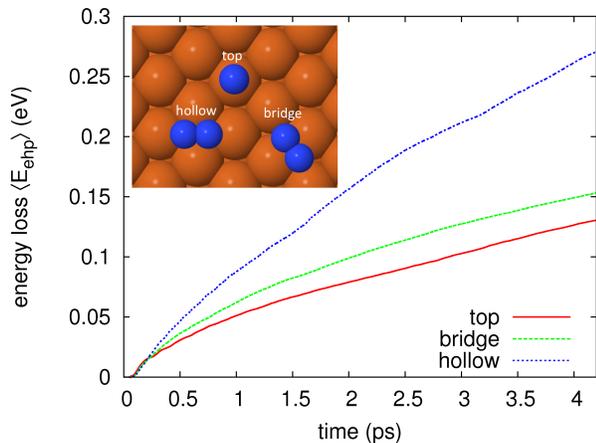


FIG. 3. Energy lost to e-h pair excitation, $\langle E_{\text{ehp}} \rangle$, during NFS+EF simulations for N_2 trajectories ending up at each adsorption site on Fe(110). Inset: molecular adsorption configurations.

TABLE I. Molecular adsorption energies E_{ads} , center of mass height Z_{CM} and molecule bondlength r at the three energy minima configurations of N_2 on Fe(110) (see Fig. 3 inset). The corresponding adsorption probabilities S_0 obtained from NFS+EF simulations are also shown.

Site	E_{ads} (eV)	Z_{CM} (Å)	r (Å)	S_0
top	0.312	2.51	1.14	0.288
bridge	0.222	1.36	1.28	0.300
hollow	0.493	1.36	1.26	0.125

the molecule (atom) to the surface lattice. In fact, this fast energy transfer fully accounts for adsorption in the $\text{N}_2/\text{Fe}(110)$ system and is its main cause in $\text{N}/\text{Ag}(111)$. After this short period of time, while $\langle E_{\text{ph}} \rangle$ reaches a plateau, E_{ehp} increases monotonically. Interestingly, we observe that, despite the phononic channel dominates N and N_2 relaxation (see Fig. 1), the energy lost to e-h pairs is by no means negligible. Note that e-h pair excitation is much more important in $\text{N}/\text{Ag}(111)$ than in $\text{N}_2/\text{Fe}(110)$, though the initial kinetic energy of the projectile is larger in the latter. This is, in fact, related to the larger adsorption energies of N on $\text{Ag}(111)$ than of N_2 on $\text{Fe}(110)$. The picture that emerges for these two systems is the following. At the first stages of the interaction the projectile transfers energy mainly to the lattice atoms in large momentum transfer collisions and gets trapped on the surface. Subsequently, it may travel as a hot species before being accommodated in the adsorption well by exciting, mainly, e-h pairs. Therefore, the amount of electronic excitation is closely related to the adsorption energy, i.e. to the energy that the adsorbate must dissipate to get fully relaxed. This explains the reported scaling of the intensity of chemicurrents measured upon adsorption of different gas species with their adsorption energy [4].

The richness of the $\text{N}_2/\text{Fe}(110)$ system regarding its adsorption properties allows us to characterize better this effect. In this system, the 6D PES calculated within the frozen surface approximation features two molecular adsorption minima at the top (local minimum) and hollow (global minimum) sites, where the N-N bond lies perpendicular and parallel to the surface, respectively [29, 31]. Remarkably, if relaxation of Fe atoms is allowed, a new molecular adsorption local minimum appears at the surface bridge site with parallel orientation (see Fig. 3 inset and Table I). The existence of such distinct adsorption wells permits to isolate the effect of E_{ads} on the e-h pair excitations from other factors. Figure 3 shows that the lowest E_{ehp} occurs for adsorption on the top site, despite its E_{ads} is about 100 meV larger than on the bridge well. The reason is that the molecules adsorbed on top are farther away from the surface, where the electron density and, therefore, the probability to excite e-h pairs is considerably smaller. The scaling of E_{ehp} with E_{ads} is recovered when comparing the results for adsorption on the hollow and bridge wells, since in both cases the molecule is close to the surface in regions of relatively high and similar density.

In summary, we have adapted the AIMDEF methodology to accurately account for the surface electron density variations caused by the surface atoms motion. With this new feature, the AIMDEF scheme is now perfectly suited to treat from first principles all possible gas-surface elementary processes where surface temperature and e-h pair excitations are determinant, such as adsorption, scattering and desorption of both light and heavy gas species. This new AIMDEF method is used here to study the role of e-h pair and phonon excitations during the thermalization of hot species on metal surfaces. Different representative systems allow us to extract general conclusions for reactions on surfaces. Thermalization of light hot reactants and intermediate products on the surface, e.g. H, will happen at a faster rate than for heavier species that involve C, N, O, because e-h pairs are more efficiently excited by the lighter atoms. For heavier atoms, we find that dissipation is dominated by lattice vibrations mainly at the initial stages of the hot species interaction with the surface, along with a substantial excitation of e-h pairs that is active during long timescales. We find that more energy is diverted into the e-h pair channel for higher adsorption energies, in consistency with the experimental observations on chemicurrents, although deviations from that behavior can be induced by the particular details of the surface electron density distribution. These conclusions can only be drawn within the theory level used here. Ultimately, the weight of each channel in the energy loss is the result of the non-trivial interplay of atomic masses, PES topography and surface electronic structure.

This work has been supported by the Basque Departamento de Educación, Universidades e Investigación, the

University of the Basque Country UPV/EHU (Grant No. IT-756-13) and the Spanish Ministerio de Economía y Competitividad (Grant No. FIS2013-48286-C02-02-P). M.B.-R. acknowledges financial support from the European Commission (Grant No. FP7-PEOPLE-2010-RG276921). The authors thankfully acknowledge the computer resources, technical expertise and assistance provided by the Red Española de Supercomputación and by the DIPC computing center.

-
- [1] G. Armand and J. R. Manson, *Phys. Rev. Lett.* **53**, 1112 (1984).
- [2] Y. Li and G. Wahnström, *Phys. Rev. Lett.* **68**, 3444 (1992).
- [3] M. Head-Gordon and J. C. Tully, *J. Chem. Phys.* **103**, 10137 (1995).
- [4] B. Gergen, H. Nienhaus, W. H. Weinberg, and E. W. McFarland, *Science* **294**, 2521 (2001).
- [5] H. Nienhaus, *Surf. Sci. Rep.* **45**, 1 (2002).
- [6] J. R. Trail, M. C. Graham, D. M. Bird, M. Persson, and S. Holloway, *Phys. Rev. Lett.* **88**, 166802 (2002).
- [7] A. C. Luntz and M. Persson, *J. Chem. Phys.* **123**, 074704 (2005).
- [8] M. Lindenblatt and E. Pehlke, *Phys. Rev. Lett.* **97**, 216101 (2006).
- [9] M. Timmer and P. Kratzer, *Phys. Rev. B* **79**, 165407 (2009).
- [10] A. Gross, *Phys. Rev. Lett.* **103**, 246101 (2009).
- [11] G. Fuchs, T. Klamroth, S. Monturet, and P. Saalfrank, *Phys. Chem. Chem. Phys.* **13**, 8659 (2011).
- [12] J. Meyer and K. Reuter, *New J. Phys.* **13**, 085010 (2011).
- [13] L. Martin-Gondre, M. Alducin, G. A. Bocan, R. Díez Muiño, and J. I. Juaristi, *Phys. Rev. Lett.* **108**, 096101 (2012).
- [14] S. M. Janke, M. Pavanello, G. J. Kroes, D. J. Auerbach, A. M. Wodtke, and A. Kandratsenka, *Z. Phys. Chem.* **227**, 1467 (2013).
- [15] M. Pavanello, D. J. Auerbach, A. M. Wodtke, M. Blanco-Rey, M. Alducin, and G. J. Kroes, *J. Phys. Chem. Lett.* **4**, 3735 (2013).
- [16] M. Blanco-Rey, J. I. Juaristi, R. Díez Muiño, H. F. Busnengo, G. J. Kroes, and M. Alducin, *Phys. Rev. Lett.* **112**, 103203 (2014).
- [17] J. Meyer and K. Reuter, *Angew. Chem. Int. Ed.* **53**, 4721 (2014).
- [18] J. T. Kindt, J. C. Tully, M. Head-Gordon, and M. A. Gomez, *J. Chem. Phys.* **109**, 3629 (1998).
- [19] J. V. Barth, *Surf. Sci. Rep.* **40**, 75 (2000).
- [20] H. Hedgeland, P. Fouquet, A. P. Jardine, G. Alexandrowicz, W. Allison, and J. Ellis, *Nat. Phys.* **5**, 561 (2009).
- [21] K. R. Harikumar, J. C. Polanyi, A. Zabet-Khosousi, P. Czekala, H. Lin, and W. A. Hofer, *Nat. Chem.* **3**, 400 (2011).
- [22] B. Schindler, D. Diesing, and E. Hasselbrink, *J. Chem. Phys.* **134**, 034705 (2011).
- [23] E. Quintas-Sánchez, C. Crespos, P. Larrégaray, J.-C. Rayez, L. Martin-Gondre, and J. Rubayo-Soneira, *J. Chem. Phys.* **138**, 024706 (2013).
- [24] P. Saalfrank, J. I. Juaristi, M. Alducin, M. Blanco-Rey, and Díez Muiño, *J. Chem. Phys.* **141**, 234702 (2014).
- [25] J. I. Juaristi, M. Alducin, R. Díez Muiño, H. F. Busnengo, and A. Salin, *Phys. Rev. Lett.* **100**, 116102 (2008).
- [26] F. L. Hirshfeld, *Theoret. Chim. Acta* **44**, 129 (1977).
- [27] S. P. Rittmeyer, J. Meyer, J. I. Juaristi, and K. Reuter, *Phys. Rev. Lett.* **115**, 046102 (2015).
- [28] L. Martin-Gondre, G. A. Bocan, M. Alducin, J. I. Juaristi, and R. Díez Muiño, *Comp. Theo. Chem.* **990**, 126 (2012).
- [29] I. Goikoetxea, M. Alducin, R. Díez Muiño, and J. I. Juaristi, *Phys. Chem. Chem. Phys.* **14**, 7471 (2012).
- [30] The Ag(111) surface is represented by 5-layer slabs separated by $d = 14.41 \text{ \AA}$ of vacuum and (2×2) lateral periodicity. For Fe(110), we use 4-layer slabs, $d = 14.15 \text{ \AA}$ and (3×3) periodic cells. We take projector augmented wave (PAW) potentials to describe the ion cores [35] and the generalized gradient approximation for the exchange and correlation functional in the PW91 [36] (RPBE [37]) formulation for N/Ag(111) ($\text{N}_2/\text{Fe}(110)$). For the plane-wave basis [38], we use an energy cut-off of 400 eV and a $5 \times 5 \times 1$ ($3 \times 3 \times 1$) Monkhorst-Pack mesh [39] for N/Ag(111) ($\text{N}_2/\text{Fe}(110)$).
- [31] I. Goikoetxea, J. I. Juaristi, R. Díez Muiño, and M. Alducin, *Phys. Rev. Lett.* **113**, 066103 (2014).
- [32] P. M. Echenique, R. M. Nieminen, J. C. Ashley, and R. H. Ritchie, *Phys. Rev. A* **33**, 897 (1986).
- [33] M. J. Puska and R. M. Nieminen, *Phys. Rev. B* **27**, 6121 (1983).
- [34] H. Winter, J. I. Juaristi, I. Nagy, A. Arnau, and P. M. Echenique, *Phys. Rev. B* **67**, 245401 (2003).
- [35] P. E. Blöchl, *Phys. Rev. B* **50**, 17953 (1994).
- [36] J. P. Perdew, J. A. Chevary, S. H. Vosko, K. A. Jackson, M. R. Pederson, D. J. Singh, and C. Fiolhais, *Phys. Rev. B* **46**, 6671 (1992).
- [37] B. Hammer, L. B. Hansen, and J. K. Nørskov, *Phys. Rev. B* **59**, 7413 (1999).
- [38] G. Kresse and J. Hafner, *Phys. Rev. B* **47**, 558 (1993).
- [39] H. J. Monkhorst and J. D. Pack, *Phys. Rev. B* **13**, 5188 (1976).